

**ANALYTICAL STUDY OF NANO MAGNETIC DRUG DELIVERY AND TARGETING IN COMPARISON WITH EXPERIMENTAL RESULTS**

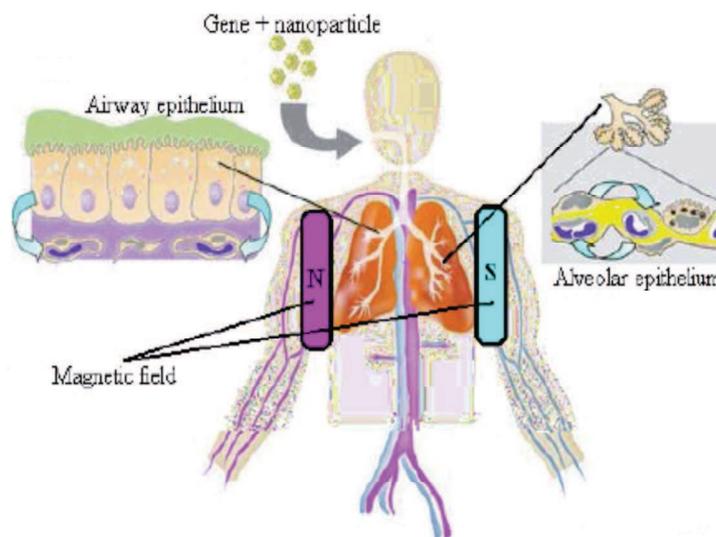
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**Abstract:** Theoretical analysis and possible application of engineered magnetic nano particles and micro particles in the lung therapy has already been studied. Several equipment for more effective delivery of drugs into the lung epithelium has been proposed earlier. This study provides an analytical method needed for development of drug targeting to lung in comparison with experimental results.

**Keywords:** Aerosolizationtherapy, Drugtargeting, magneticnanoparticle, Quadrapole magnet.

**Introduction:** One of the basic requirements in drug management is that sufficient quantities of the drug reach the desired site in the body. The main aim is to decrease the amount of drug to the non target sites as it can harm the healthy organs and cells of the body. This is achieved by attaching the drugs to the magnetic nano particles which act as carrier particles and are controlled externally by magnetic fields. This is done by locating the permanent magnets directly over the affected sites in the body.

Aerosolization therapy for lung diseases has been applied for a long time. The lung is a perfect point of entry not only for local therapy but also for systematic delivery of therapeutic substances. Recent technical improvements, with regard to both the production and delivery of aerosols indicate that appropriate inhalation techniques may be the therapeutic way forward for a variety of diseases. Since the nano carrier system may be administered to the airways easily, a number of respiratory diseases may be approached using nano particles.



**Principle of magnetic drug targeting to the lungs:**

**Methods:**

**Theoretical Model:** There are several forces acting on magnetic particle in viscous environment and magnetic field, such as magnetic force due to all field sources, Stokes’ viscous drag force, inertia, buoyancy and gravity, thermal kinetics (Brownian motion) particle fluid interactions and interparticle effects such as magnetic dipole interactions, electric double layer interactions, and van der Walls force. Except magnetic, inertial, and viscous drag force the other interactions are negligible for magnetic micro- or nanoparticles. For this reason, we suppose in this

paper only these dominant effects. The trajectory of motion of magnetic particle in magnetic field and viscous fluid ambient using Newton’s law:

$$m \frac{dv}{dt} = F_m + F_s(\mathbf{1})$$

where  $m$  and  $v$  are the mass and velocity of the particle, and  $F_m$  and  $F_s$  are the magnetic and stokes drag force respectively.

**Magnetic Forces:** A magnetic particle is replaced by “equivalent point dipole moment  $\mathbf{m}_{p,eff}$  localized at the center of particle. According to this method, force acting on the dipole is given by

$$F_m = \mu(\mathbf{m}_{p,eff} \cdot \nabla) H_a \tag{2}$$

where  $\mu$  is permeability of fluid ambient,  $\mathbf{m}_{p,eff}$  is

“effective” dipole moment of the particle, which depends on externally applied magnetic field intensity  $H_a$  at the center of particle, where the dipole moment is localized. A magnetic force is therefore a function of external magnetic field gradient and the magnetization of the particle. Below the saturation the particles are linearly magnetized with their magnitude of magnetic moment increasing in the direction of external field. Beyond the saturation point, magnetic moment magnitude tends to a constant value. According to this magnetization model of particles based on self demagnetization and magnetic saturation developed by Furlani group the effective dipole moment can be expressed as

$$m_{p,eff} = V_p f(H_a) H_a \tag{3}$$

We consider magnetic particle with radius  $R_p$  and volume  $V_p$  and a function

$$f(H_a) = \frac{3(\chi_p - \chi_f)}{(\chi_p - \chi_f) + 3} H_a < \frac{3\chi_p}{3\chi_p} M_{sp} \tag{4}$$

$$= \frac{M_{sp}}{H_a} H_a \geq \frac{3\chi_p}{3\chi_p} M_{sp}$$

Where  $\chi_p$  and  $\chi_f$  are the magnetic susceptibilities of the particle and ambient fluid, respectively,  $M_{sp}$  is saturation magnetization of the particle, and  $H_a = |\mathbf{H}_a|$ . We assume nonmagnetic fluid ( $\chi_f \approx 0$ ) and high susceptibility of magnetic particles, i.e.  $\chi_p \gg 1$ , which is in the case of water or air as fluid ambient, and magnetite  $Fe_3O_4$  as particles accomplished; hence

$$B/\mu < M_{sp} / 3 \tag{5}$$

where  $B$  is magnetic flux density of external field and is valid:  $B/\mu = H_a$ .

As an illustration we could calculate trajectory of magnetite particles in the magnetic field of permanent quadrupole. Magnetic flux density was modelled as magnetostatic problem in which the fields are time-independent. In this case, the field intensity ( $H$ ) and flux density ( $B$ ) must obey equations

$$\nabla \times \mathbf{H} = \mathbf{J},$$

$$\nabla \cdot \mathbf{B} = 0$$

with a constitutive relationship between  $B$  and  $H$  for each material

$$\mathbf{B} = \mu \mathbf{H}.$$

If a material is nonlinear (e.g. saturating iron or alnico magnets), the permeability,  $\mu$  is actually a function of  $B$ :  $\mu = B/H(B)$ .

**Quadrupole**

**Mathematical model:** A planar model which is not 3D, but it represents the distribution of flux density in transversal plane of infinite long permanent magnets arranged to quadrupole with zero value of perpendicular component of flux density was created.

The quadrupole consists of eight sphenoid blocks, in the section, of uniformly magnetized Neodymium rare earth magnets with magnetic energy product 37 MG.Oe (megagauss-oersted; NdFeB N37) and with magnetization orientation revolved in 135° between adjacent blocks. The geometry of quadrupole is determined by the radii of inscribed and circumscribed circle, i.e. 1.4 and 4.5 cm in our case, respectively. The purposes of magnetic targeting in a long, is illustrated in the Fig. In this model we consider in addition to magnetic force, fluidic force acting on a moving particle in fluid medium. Its magnitude is determined by Stokes’ law for the drag on a sphere with radius  $R_p$  in uniform flow,

$$F_s = -6\pi\eta R_p (v_p - v_f) \tag{6}$$

where  $\eta$  and  $v_f$  are the viscosity and the velocity of the fluid, respectively, and  $v_p$  is the velocity of the particle. In this case is the fluid ambient quiescent, i.e.  $v_f = 0 \text{ m. s}^{-1}$ .

**Calculating Flux Density Components Of A Quadrupole:** In the next step magnetic flux density of permanent quadrupole is computed by determining the  $B$ -field outside an infinitely long cylindrical magnet that has an alternating radial polarization and a linear second quadrant demagnetization curve of the form

$$\mathbf{B} = \mu \mathbf{H} + \mu_0 \mathbf{M}_s(\varphi)$$

Let  $R_1$  and  $R_2$  denote the inner and outer radii of the cylinder, respectively. As the magnet is infinitely long, the problem reduces to a 2D boundary value problem (BVP). We use the solution method presented in [2] with the magnetization inside the cylinder given by

$$M_s(\varphi) = M(\varphi) \frac{1}{r}$$

$$M(\varphi) = \sum_{i=1}^{\infty} M_i \cos(i\varphi)$$

$$M_i = \frac{M_s}{i\pi} \sum_{k=1}^{N_{pole}} N_{pole} \cos$$

$$(-1)^k [\sin^{-1}((2k-1)\pi/N_{pole}) - \sin^{-1}(2k+1)\pi/N_{pole}]$$

$$M_i = ]$$

$N_{pole}$  is the number of poles

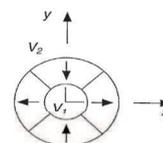


FIGURE 4.18 Infinite cylinder with radial magnetization.

A solution is sought for the region  $V_2$  outside the cylinder. The coefficients of interest are  $V_i^{(2)}$  and  $U_i^{(2)}$ . From the Orthogonality of  $\sin(i\varphi)$  and  $\cos$

(iφ) we find that  $V_i^{(2)} = 0$  for all i. The coefficients  $U_i^{(2)}$  are given by  $U_i^{(2)}(M_i, R_1, R_2, \mu)$  which is calculated numerically as per [2].

$$B_r^{(2)}(r, \phi) = \mu_0 \sum_{i=1}^{\infty} i r^{-(i+1)} U_i^{(2)} \cos(i\phi)$$

$$B_\phi^{(2)}(r, \phi) = \mu_0 \sum_{i=1}^{\infty} i r^{-(i+1)} U_i^{(2)} \sin(i\phi) \quad (7)$$

are calculated. Movement of magnetic particles in the plane in magnetic field with flux density  $B$  in fluid ambient with viscosity  $\eta$ , which is not moving, is described by system of **ordinary differential equations (ODE)**:

$$dx/dt = v_{p,x}$$

$$dy/dt = v_{p,y}$$

$$dv_{p,x}/dt = \frac{1}{mp} \left\{ \frac{V_p}{\mu} f(B(x,y)) \left[ B_x(x,y) \frac{\partial B_x(x,y)}{\partial x} + B_y(x,y) \frac{\partial B_x(x,y)}{\partial y} \right] - 6\pi\eta R_p v_{p,x} \right\}$$

$$dv_{p,y}/dt = \frac{1}{mp} \left\{ \frac{V_p}{\mu} f(B(x,y)) \left[ B_x(x,y) \frac{\partial B_y(x,y)}{\partial x} + B_y(x,y) \frac{\partial B_y(x,y)}{\partial y} \right] - 6\pi\eta R_p v_{p,y} \right\}$$

(8)

where  $mp = V_p \rho_p$  and  $V_p = 4/3\pi r^3$  are the mass and volume of the particles respectively

This system of equations is obtained by combinations of expressions from (1) to (6)

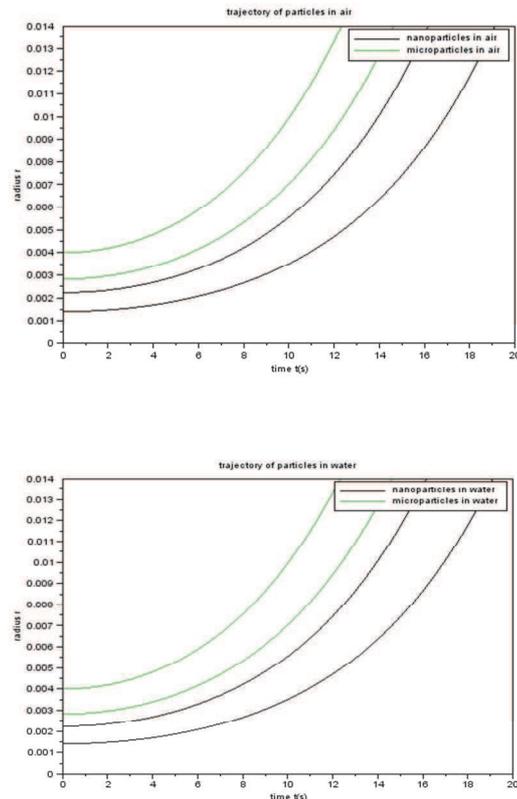
Gradient of flux density components is calculated using equation (7).

The field is expressed as  $B_0 = (B_r^2 + B_\phi^2)^{1/2}$  where  $B_0$  is both the field coefficient and gradient of the Quadrupole. Setting the trajectory of magnetite ( $Fe_3O_4$ ) particles, with density  $\rho = 5000 \text{ kg.m}^{-3}$  and a saturation magnetization  $M_{sp} = 4.78 \times 10^5 \text{ A.m}^{-1}$ , in the

magnetic field created by quadrupole consisting of permanent NdFeB  $N_{37}$  magnets, which is calculated analytically using eqn(7) in nonmoving and nonmagnetic fluid, with viscosity equal to that of water or air, i.e.  $\eta = 1.003 \times 10^{-3} \text{ N.s.m}^{-2}$  or  $\eta = 1.82 \times 10^{-5} \text{ N.s.m}^{-2}$ , respectively. For this particle  $B_0$  was found to be  $0.018 \text{ T}$ .

**Computations Using Scilab:** Scilab 5.4 numerical solver was used for solving the set of equations (8). Computations were done for nanoparticles with radius  $50 \text{ nm}$ , which are often used in magnetic drug targeting, as well as microparticles with radius  $10 \mu\text{m}$ , which can be used for magnetic separation. The trajectories of magnetite particles of each size in the air as fluid ambient as shown by [2] was solved analytically and the trajectory traced.

Initial conditions for calculations were: randomly position in the circle with radius  $0.4 \text{ cm}$  from the center of quadrupole and zero initial velocity. Computation of trajectory for each particle was stopped after it reached internal boundary of quadrupole, approximately at a distance  $1.4 \text{ cm}$  from the center. We can see microparticles were captured faster than nanoparticles. Also small nanoparticles were attracted in a longer, but still reasonable time approximately after  $16 \text{ s}$ . Trajectories of magnetite particles with different sizes, microparticles with radius  $10 \mu\text{m}$  and nanoparticles with radius  $50 \text{ nm}$ , in the field of permanent quadrupole in the air and in water in different times are traced using scilab. Mean capture time is in this case shorter than ones in viscously water medium, due to smaller drag force. Capture time of magnetic particles in magnetic field of quadrupole, can be also decreased by increasing magnetic flux densities and field gradients of studied configuration replacing permanent magnets by pulsed electromagnets, similar to those pulsed quadrupole lenses using for focusing heavy ion beams in accelerator physics. These results equally apply also to magnetic separation of proteins, DNA, and whole cells.



**Conclusions:** The trajectories of magnetite particles of different sizes in the field of permanent quadrupole in the air and water were traced in different times with the help of numerical solver. The tendency of the particle to be captured by the magnet increases when air is chosen as the medium. The analytical model presented here was compared with the

experimental model [3] and found to be correct. This model presented here is ideal for parametric analysis of magnetic targeting *in vivo*. The numerical procedure proved to be a powerful tool for further development of the magnetic drug targeting technique for patient specific clinical studies.

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