

**ON VARIOUS DECOMPOSITIONS OF A SIGNED LATTICE MEASURE**

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**Abstract:** This paper describes that the concepts of lattice measurable space, lattice measure space, lattice finite measure and lattice finite measure space and the perceptions of lattice  $\sigma$  - finite measure, lattice  $\sigma$  - finite measure space, absolutely continuous lattice measure and also establishes the various decompositions of a signed lattice measures.

**Keywords:** lattice measurable space, lattice  $\sigma$  - finite measure space, absolutely continuous lattice measure and signed lattice measures.

**Introduction:** In 2005 Mona Khare and BhawnaSingh [2] introduced the notation of a signed measure arises if a measure is allowed to take both positive and negative values. A set that is both positive and negative with respect to a signed measure is termed as a null set. They generalized some concepts in measure theory by means of class of null set. In 1963, GabourSzasz [3] Introduction To Lattice Theory explained much about sublattice concept like if  $P(M)$  denotes the set of all sub sets of  $M$  including the void set is a lattice with respect to the operations of set union and intersection. This lattice is called sublattice of  $M$ . He also explained that the set union and the set intersection of any two closed sets of a plane are likewise closed. Hence, The closed sets of a plane form a lattice with respect to these set operations. For similar reasons, the corresponding statement holds for all open sets of a plane also every single element sub set of a partially ordered set is an interval and this interval is a sub lattice, also this single element sub set  $\{a\}$  of a lattice  $L$  can be recorded as a zero length chain or as an interval  $[a, a]$ , so every  $\{a\}$  is a sublattice of  $L$  as well as measurable  $[m(a) = 0]$ .

In section2, we defined the concepts of lattice measurable space, lattice measure space, lattice finite measure and lattice finite measure space and the perceptions of lattice  $\sigma$  - finite measure, lattice  $\sigma$  - finite measure space, absolutely continuous lattice measure and also establish the various decompositions of a signed lattice measures.

**Various Decompositions Of A Signed Lattice Measure:**

**Definition2.0.** Let  $X$  is the entire set and  $L$  is a lattice of any subsets of  $X$ . If a lattice  $L$  satisfies the following conditions, then it is called a lattice  $\sigma$ -Algebra;

1.  $\forall h \in L, h^c \in L.$
2. If  $h_n \in L$  for  $n = 1, 2, 3, \dots$ , then  $\bigcup_{n=1}^{\infty} h_n \in L.$

We denote  $\beta$ , as the lattice  $\sigma$ -Algebra generated by  $L$ .

**Definition2.1.** By a signed lattice measure on the lattice measurable space  $(X, \beta)$  we mean a mapping  $\nu: \beta \rightarrow \mathbb{R} \cup \{\infty\}$  or  $\mathbb{R} \cup \{-\infty\}$ , satisfying the following properties

- (1)  $\nu(\emptyset) = \nu(0) = 0.$
- (2)(i) For all  $h, g \in \beta$  such that

- $\nu(h), \nu(g) \geq 0; h \leq g \Rightarrow \nu(h) \leq \nu(g).$
- (ii) For all  $h, g \in \beta$  such that  $\nu(h), \nu(g) \leq 0; h \leq g \Rightarrow \nu(h) \leq \nu(g).$
- (3) For all  $h, g \in \beta,$   
 $\nu(h \vee g) + \nu(h \wedge g) = \nu(h) + \nu(g).$
- (4) If  $h_n \in \beta, n \in \mathbb{N}$  such that  $h_1 \leq h_2 \leq \dots \leq h_n \leq \dots$ , then  
 $\nu(\bigcup_{n=1}^{\infty} h_n) = \lim \nu(h_n).$

**Definition2.2.** The ordered pair  $(X, \beta)$  is said to be lattice measurable space.

**Definition2.3.**  $A$  is called a positive lattice if  $A$  is lattice measurable and for any lattice measurable set  $E$  in  $A, \nu(E) \geq 0$ ; Similarly,  $B$  is called a negative lattice if  $B$  is a lattice measurable and for any lattice measurable set  $E$  in  $B, \nu(E) \leq 0.$

**Definition2.4. (Positive and Negative parts of  $\nu$ ):** Let  $(X, \beta)$  be a lattice measurable space and  $\nu$  be a signed lattice measure defined on  $\beta$ . If we define  $\nu^+$  and  $\nu^-$  on for all  $E \in \beta$  such that  $\nu^+(E) = \nu(E \wedge A)$  and  $\nu^-(E) = -\nu(E \wedge B)$  with  $\nu = \nu^+ - \nu^-$ . Then  $\nu^+$  and  $\nu^-$  are called the positive and the negative parts of a signed lattice measure respectively. Where  $\{A, B\}$  be a lattice Hann decomposition on  $X$ .

**Definition2.5. (Mutually singular lattice measures):** Let  $(X, \beta)$  be a lattice measurable space and  $\nu$  be a signed lattice measure defined on  $\beta$ . If we define two measures  $\nu_1$  and  $\nu_2$  are said to be mutually singular lattice measures with each other to be denoted by  $\nu_1 \perp \nu_2$ . That is if there exists two lattice measurable sets  $A$  and  $B$  such that  $A \vee B = X$  and  $A \wedge B = \emptyset$  and  $\nu_1(A) = \nu_2(B) = 0.$

**Definition2.6.** The lattice measurable space  $(X, \beta)$  together with a lattice measure  $\mu$  is called a lattice measure space and it is denoted by  $(X, \beta, \mu).$

**Definition2.7.** Let  $(X, \beta, \mu)$  be a lattice measure space. If  $\mu(X)$  is finite, then  $\mu$  is called lattice finite measure.

**Definition2.8.** If  $\mu$  is a lattice finite measure then  $(X, \beta, \mu)$  is called a lattice finite measure space.

**Definition2.9.** Let  $(X, \beta, \mu)$  be a lattice measure space. If there exists a sequence of lattices measurable sets  $\{X_n\}$  such that

$X = \bigvee_{n=1}^{\infty} X_n$  (ii)  $\mu(X_n)$  is finite.

Then  $\mu$  is called a lattice  $\sigma$ -finite measure.

**Definition 2.10.** If  $\mu$  be a lattice  $\sigma$ -finite measure, then  $(X, \beta, \mu)$  is called lattice  $\sigma$ -finite measure space.

**Definition 2.11. Absolutely continuous lattice measure:** Let  $(X, \beta, \mu)$  be a lattice measure space and  $\nu$  be a lattice measure defined on  $\beta$  for each sublattice  $A$  of  $X$ , whenever  $\mu(A) = 0$  implies  $\nu(A) = 0$ . Then  $\nu$  is said to be absolutely continuous lattice measure with respect to  $\mu$  and it is denoted by  $\nu \ll \mu$ .

By using the above concepts one can establish the following theorems.

**Theorem 2.1. :**

**Lattice Hahn-Decomposition:** Let  $(X, \beta)$  be a lattice measurable space and  $\nu$  be a signed lattice measure defined on  $\beta$ . Then there exists a positive lattice  $A$  and a negative lattice  $B$  such that  $A \vee B = X$  and  $A \wedge B = \emptyset$ .

**Theorem 2.2.**

**Jordan-Decomposition On Signed Lattice Measure:** Let  $(X, \beta)$  be a lattice measurable space and  $\nu$  be a signed lattice measure defined on  $\beta$ . Then there exists a unique pair of mutually singular lattice measures  $\nu^+$  and  $\nu^-$  with  $\nu = \nu^+ - \nu^-$ .

**Theorem 2.3. Radon - Nikodym Theorem On Signed Lattice Measure:** Suppose  $(X, \beta, \mu)$  be a lattice  $\sigma$ -finite measure space and  $\nu$  be a lattice measure defined on  $\beta$  which is absolutely continuous with respect to  $\mu$ . Then there exists, a non-negative lattice measurable function  $f$  such that for all  $E$  in  $\beta$  we have  $\nu(E) = \int_E f d\mu$ . More

over the lattice measure  $f$  is unique in this sense that if  $g$  is also a non-negative lattice measurable function such that  $\nu(E) = \int_E g d\mu, E \in \beta$  then  $f = g$  almost everywhere on  $X$  with respect to  $\mu$ .

**Theorem 2.4.**

**Lebesgue Decomposition Of A Signed Lattice Measure:**

Suppose  $(X, \beta, \mu)$  be a lattice  $\sigma$ -finite measure space and  $\nu$  be a lattice  $\sigma$ -finite measure defined on  $\beta$ . Then we find

- (1) A lattice measure  $\nu_0$  mutually singular lattice measure with respect to  $\mu$  and
- (2) A lattice measure  $\nu_1$  absolutely continuous lattice measure with respect to  $\mu$  such that
- (3)  $\nu = \nu_0 + \nu_1$  more over  $\nu_0$  and  $\nu_1$  are unique.

**Proof.** Let  $(X, \beta, \mu)$  be a lattice  $\sigma$ -finite measure space and  $\nu$  be a lattice  $\sigma$ -finite measure defined on  $\beta$ . Now  $\mu$  and  $\nu$  are lattice  $\sigma$ -finite measures, therefore  $\lambda = \mu + \nu$  is also a lattice  $\sigma$ -finite measure. Also evidently,  $\mu \ll \lambda$  and  $\nu \ll \lambda$ . Therefore by (Radan - Nikodym Theorem) there exists two non-negative lattice measurable functions  $f$  and  $g$  on  $X$  such that for every  $E \in \beta$ ,

We have  $\mu(E) = \int_E f d\lambda$  (1) and

$$\nu(E) = \int_E g d\lambda \quad (2)$$

Now  $f$  is a non-negative lattice measurable function defined on  $X$ . Define two sets  $A$  and  $B$  such that  $A = \{x \in X / f(x) > 0\}$  and  $B = \{x \in X / f(x) = 0\}$ . Clearly  $A \vee B = X$  and  $A \wedge B = \emptyset$

from (1)  $\mu(B) = \int_B f d\lambda = 0$ . Therefore  $\mu(B) = 0$

Case (1). Define a measure  $\nu_0$  such for every  $E \in \beta$ , we have  $\nu_0(E) = \nu(E \wedge B)$ .

Now  $\nu_0(A) = \nu(A \wedge B) = \nu(\emptyset)$ . Therefore there exists two lattice measurable sets  $A$  and  $B$  such that  $A \vee B = X$  and  $A \wedge B = \emptyset$  also  $\nu_0(A) = \nu_0(B) = 0$ . Therefore  $\nu_0$  and  $\mu$  are mutually singular lattice measures.

Case (2). Define a lattice measure  $\nu_1$  such for every  $E \in \beta$ , we have  $\nu_1(E) = \nu(E \wedge A)$ .

Let  $E$  be a lattice measurable set of  $\mu$  measure zero.

Now  $\mu(E) = 0$ . Implies  $\int f d\lambda = 0$  (since (1))

Implies  $\lambda(E) = 0$ . Implies  $\lambda(E \wedge A) = 0$ .

Implies  $\mu(E \wedge A) + \nu(E \wedge A) = 0$  (since  $\lambda = \mu + \nu$ )

Implies  $0 + \nu(E \wedge A) = 0$  (since  $\mu(E) = 0$ ).

Implies  $\nu(E \wedge A) = 0 = \nu_1(E) = 0$ .

Therefore  $\mu(E) = \nu_1(E) = 0$ . Hence  $\nu_1$  absolutely continuous lattice measure with respect to  $\mu$

Case (3). For every  $E \in \beta$ , we have

$$E = (E \wedge A) \vee (E \wedge B)$$

Implies  $\nu(E) = \nu(E \wedge A) + \nu(E \wedge B)$

Implies  $\nu(E) = \nu_0(E) + \nu_1(E)$ . Implies  $\nu = \nu_0 + \nu_1$ .

**Uniqueness:** Suppose there exists two pair of lattice measures  $\nu_0, \nu_1$  and  $\nu_0^1, \nu_1^1$

Such that  $\nu = \nu_0 + \nu_1$  (3)

and  $\nu = \nu_0^1 + \nu_1^1$  (4).

Let  $\{A, B\}$  and  $\{A^1, B^1\}$  be two pair of lattice measurable sets as explained above and

$$A \vee B = X, A \wedge B = \emptyset \text{ and } A^1 \vee B^1 = X, A^1 \wedge B^1 = \emptyset$$

$$\mu(B) = \mu(B^1) = \nu_0(A) = \nu_0(A^1) = 0 \quad (5)$$

For every  $E \in \beta$ , we have  $E = (E \wedge A \wedge A^1) \vee (E \wedge A \wedge B) \vee (E \wedge A^1 \wedge B) \vee (E \wedge B \wedge B^1) \dots$  (6).

Since  $\mu(B) = \mu(B^1) = 0$ . The  $\mu$  lattice measure of last three terms of (6) are zero. But  $\nu_1 \ll \mu$

Implies the  $\nu_1$  lattice measure of last three terms of (6) are zero. (3) - (4) gives  $(\nu_0 - \nu_0^1) + (\nu_1 - \nu_1^1) = 0$ . Implies  $(\nu_0 - \nu_0^1) = (\nu_1^1 - \nu_1)$  (7). In view of the above argument for every  $E \in \beta$ ,

$$(\nu_1^1 - \nu_1)(E) = (\nu_1^1 - \nu_1)(E \wedge A \wedge A^1). \text{ Implies } (\nu_1^1 - \nu_1)(E) = (\nu_0 - \nu_0^1)(E \wedge A \wedge A^1) = 0 \text{ (since (5))}$$

This is true for every  $E \in \beta$ , that is  $\nu_1^1 - \nu_1 = 0$ . Implies  $\nu_1^1 = \nu_1$ . Therefore  $\nu_1$  is unique lattice measures.

Now from (7)  $\nu_0 - \nu_0^1 = 0$ . Implies  $\nu_0 = \nu_0^1$ . Implies  $\nu_0$  is unique lattice measure.

Therefore there exists a unique pair of lattice measures  $\nu_0$  and  $\nu_1$  such  $\nu = \nu_0 + \nu_1$ .

**Conclusion:** In this paper defined the concepts of lattice measurable space, lattice measure space, lattice finite measure and lattice finite measure space and the perceptions of lattice  $\sigma$  – finite measure, lattice finite measure space, absolutely continuous lattice measure and also established the various decompositions of a signed lattice measures.

**References:**

1. Birkhoff. G, Lattice Theory 3rd ed., AMS Colloquium Publications, Providence, RI, 1967.
2. Mona Khare and Bhawna Singh, Weakly tight functions and their decomposition, IJMMS, 2005:18(2005)2991-2998.
3. Szasz Gabor, Introduction to lattice theory, academic press, New York and London, 1963.

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